



SYNTHETIC RESPONSE AND RADIATION ABSORPTION IMPACTS ON UNSTEADY MHD FREE CONVECTIVE FLOW OVER A VERTICAL PERMEABLE PLATE

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ABSTRACT

This article manages the impacts of substance response and radiation ingestion on insecure MHD free convection stream of a thick, incompressible electrically leading liquid along a semi-limitless vertical plate within the sight of a uniform transverse attractive field, warm and fixation buoyancy impacts. The representing nonlinear halfway differential conditions have been lessened to the coupled nonlinear common differential conditions by irritation system. Numerical evaluation of the explanatory results are performed through diagrams.

Key words: Radiation absorption, Chemical reaction, MHD, Heat transfer, Mass transfer.

INTRODUCTION

Characteristic convection stream over vertical plate drenched in permeable media has central significance in view of its potential applications in soil material science, geohydrology, and filtration of solids from fluids, substance designing and natural frameworks. Magneto convection assumes an imperative part in horticulture, petroleum commercial enterprises, geophysics and in astronomy. Radiative warmth and mass exchange assume an essential part in assembling commercial enterprises for the outline of blades, steel rolling, atomic force plants, gas turbines.

Joined warmth and mass exchange from various procedures with permeable media has a wide range engineering and industrial applications, like upgraded oil recuperation,

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underground vitality transport, geothermal stores, cooling of atomic reactors, drying of permeable solids, pressed bed synergist reactors and warm protection. Magnetohydrodynamics (MHD) is the branch of continuum mechanics, which manages the stream of electrically directing liquids in electric and attractive fields. As of late, advance has been significantly made in the investigation of warmth and mass move in magnetohydrodynamic (MHD) streams because of its application in numerous gadgets, similar to the MHD power generators and Hall quickening agents.

The developing requirement for compound responses in substance and hydro-metallurgical commercial ventures requires concentrating on warmth and mass exchange with concoction response. The nearness of a remote mass in water or air causes some sort of synthetic response. This might be available either independent from anyone else or as a blend with air or water. In numerous synthetic building forms, a compound response happens between a remote mass and the liquid in which the plate is moving. These procedures occur in various modern applications like polymer generation, producing pottery or dishes and sustenance handling. The investigation of warmth era or ingestion in moving liquids is critical in issues managing concoction responses and those worried with separating liquids.

Kafousias and Georgantopoulos¹ considered the transverse attractive consequences for the free convective stream of an incompressible, electrically directing liquid past a non-leading and non-attractive, vertical restricting surface, the overseeing conditions were settled by the typical Laplace change procedure. Raptis and Soundalgekar² decided the impacts of mass exchange on the stream of an electrically directing liquid past a relentlessly moving endless vertical permeable plate under the activity of a transverse attractive field. Raptis and Soundalgekar³ considered the issue of stream of an electrically leading liquid past a consistently moving vertical interminable plate in nearness of steady warmth flux and consistent suction at the plate and incited attractive field is likewise considered. England and Emery⁴ have considered the radiation impacts of an optically thin dark gas limited by a stationary plate. Raptis and Massalas⁵ researched the impacts of radiation on the oscillatory stream of a dim gas, engrossing transmitting in nearness incited attractive field and systematic arrangements were gotten with help of bother strategy. Mansour et al.⁶ explored the impacts of synthetic response, warm stratification, Soret and Dufour numbers on MHD free convective warmth and mass exchange of a gooey, incompressible and electrically directing liquid on a vertical extending surface installed in an immersed permeable medium. Gupta⁷ considered precarious magneto-convection under lightness powers. Chamkha⁸ has broke down the shaky MHD free three-dimensional convection over a slanted penetrable surface with warmth era/absorption. Radiation magnetohydrodynamic convection streams are likewise critical in astrophysical and geophysical administrations. Hossain et al.⁹

considered free convection-radiation interaction, which is from an isothermal inclined plate. Distinctive explores have been sent to investigate the impacts of warm radiation on various streams¹⁰⁻¹⁴.

Ramprasad et al.¹⁵ reported heat absorption for a past an inclined moving surface unsteady MHD convective heat and mass transfer. Ramana Reddy et al.¹⁶ investigated chemical response and heat radiation with MHD blended convection oscillatory over a vertical surface.

The scattering has building applications, for example, a sensible warmth stockpiles beds, fired handling, oil repository and petroleum recuperation, and so on. The warm and solutal scattering impacts turn out to be more vital when the inertial impacts are predominant. The investigation of convection through permeable media with variable consistency and warm conductivity are imperative in a few building applications, like glass fiber, drawing of plastic movies, investigation of spilling poison unrefined petroleum over the surface of the seawater, cooling of atomic reactors, nourishment handling, petroleum store operations, throwing and welding in assembling forms, and so forth. A few investigators have considered the impact of variable properties in blended convection streams.

The stream warmth and mass transport in a fluid soaked permeable medium are regularly happened in numerous designing procedures, common habitats and geophysical applications, for example, leakage of water through waterway bed, relocation of poisons into the dirt and aquifers and stream of dampness through permeable modern materials and so forth.

A precise expectation of warmth exchange coefficient has massive critical in the investigation of the softening rate of a strong body encompassed by warm liquid, for example, floating of the chunk of ice in ocean water, while it decides the rate of vitality trade. Coupled warmth and mass exchange joined by dissolving has increased more significance because of its incomprehensible number of uses lately. Just to give some examples are, softening of permafrost, silicon wafer process, throwing and welding process and so forth. The softening impact is essential in either hot expulsion or hot working procedures, as it is done over the materials re-crystallization temperature to keep materials from work solidifying and to make it simpler to push the materials through die. Singh and Gupta¹⁷ have examined the MHD free convective stream of a gooey liquid through a permeable medium limited by a wavering permeable plate in slip stream administration with mass exchange.

The investigation of warmth and mass exchange issues with synthetic response is of extraordinary pragmatic significance to specialists and researchers due to their practically

general event in numerous branches of science and building. A couple of delegate fields of enthusiasm for which joined warmth and mass exchange alongside compound response assume an imperative part are concoction process commercial enterprises, like sustenance preparing and polymer production. The impelled attractive field is disregarded by accepting a little attractive Reynolds number¹⁹⁻²¹.

Limit layer stream on nonstop moving surfaces discovers applications in various designing procedures. Expulsion of plastic or elastic sheets, moving of string between a food roll and a wind-up roll take after the investigation of nonstop moving surfaces are a few illustrations. Some precise and numerical answers for the warmth exchange issue are found by Nigam and Singh¹⁸.

The point of the present article is to explore the impact of different parameters like parameter for warmth assimilation coefficient, concoction response parameter, warm Grashof number, mass Grashof number and so on convective warmth exchange along a slanted moving vertical plate in permeable medium. The overseeing non-direct incomplete differential conditions are initially changed into a dimensionless structure and in this way coming about non-comparative arrangement of conditions has been comprehended utilizing the irritation method.

Issue solution

Consider the shaky two dimensional MHD free convective stream of a thick incompressible, electrically leading and emanating liquid in an optically thin environment past an unbounded warmed vertical permeable plate implanted in a permeable medium in nearness of warm and focus lightness impacts. Give the pivot a chance to be taken in vertically upward course along the plate and hub is typical to the plate. It is accepted that there exist a homogeneous substance response of first request with consistent rate between the diffusing species and the liquid. A uniform attractive field is connected in the bearing opposite to the plate. The gooey dissemination and the Joule warming impacts are thought to be irrelevant in the vitality condition. The transverse connected attractive field and attractive Reynolds number are thought to be little, so that the incited attractive field is irrelevant. Likewise, it is accepted that there is no connected voltage, so that the electric field is missing. The grouping of the diffusing species in the parallel blend is thought to be little in correlation with the other substance species, which are available, and subsequently the Soret and Dufour impacts are unimportant. Under the above suppositions and also Boussinesq's estimation, the conditions of protection of mass, force, vitality and focus administering the free convection limit layer stream over a vertical permeable plate in permeable medium can be communicated as:

$$\frac{\partial v'}{\partial y'} = 0 \quad \dots(1)$$

$$\frac{\partial u'}{\partial t'} + v' \frac{\partial u'}{\partial y'} = \nu \frac{\partial^2 u'}{\partial y'^2} + g\beta(T' - T'_\infty) + g\beta^*(C' - C'_\infty) - \frac{\sigma B_0^2}{\rho} u' - \frac{\nu}{k'} u' \quad \dots(2)$$

$$\frac{\partial T'}{\partial t'} + v' \frac{\partial T'}{\partial y'} = \frac{k}{\rho c_p} \frac{\partial^2 T'}{\partial y'^2} - \frac{1}{\rho c_p} Q_r (C' - C'_\infty) \quad \dots(3)$$

$$\frac{\partial C'}{\partial t'} + v' \frac{\partial C'}{\partial y'} = D \frac{\partial^2 C'}{\partial y'^2} - K_r (C' - C'_\infty) \quad \dots(4)$$

where u' , v' are the velocity components in x' , y' directions respectively, t' – the time, p' – the pressure, ρ – the fluid density, g – the acceleration due to gravity, β and β^* – the thermal and concentration expansion coefficients respectively, K' – the permeability of the porous medium, T' – the temperature of the fluid in the boundary layer, ν – the kinematic viscosity, σ – the electrical conductivity of the fluid, T'_∞ – the temperature of the fluid far away from the plate, C' – the species concentration in the boundary layer, C'_∞ – the species concentration in the fluid far away from the plate, B_0 – the magnetic induction, α – the fluid thermal diffusivity, c_p – specific heat at constant, D – the coefficient of chemical molecular diffusivity, K_r – the chemical reaction, Q_r is coefficient of heat absorption.

The boundary conditions for the velocity, temperature, and concentration fields are given as follows:

$$\begin{aligned} u' = L' \left(\frac{\partial u'}{\partial y'} \right), T' = T'_w + \varepsilon (T'_w - T'_\infty) e^{n't'}, C' = C'_w + \varepsilon (C'_w - C'_\infty) e^{n't'} & \quad \text{at } y' = 0 \\ u' \rightarrow 0, \quad T' \rightarrow T'_\infty, \quad C' \rightarrow C'_\infty & \quad \text{as } y' \rightarrow \infty \dots(5) \end{aligned}$$

Where T'_w and T'_∞ are the temperature at the wall and infinity. C'_w and C'_∞ are the species concentration at the wall and at infinity, respectively.

From Equation (1), plainly suction speed at the plate is either a constant or capacity of time as it were. Consequently the suction speed typical to the plate is thought to be in the structure –

$$v' = -V_0 \left(1 + \varepsilon A \exp^{i\omega t'} \right) \quad \dots(6)$$

Where A will be a genuine positive constant, and $\varepsilon \ll 1$, $\varepsilon A \ll 1$ is little and V_0 is a non-zero positive steady, the negative sign shows that the suction is towards the plate.

With a specific end goal to compose the administering conditions and the limit conditions in dimensionless structure, the accompanying non-dimensional amounts are presented.

$$\begin{aligned}
 u &= \frac{u'}{V_0}, v = \frac{v'}{V_0}, y = \frac{V_0 y'}{\nu}, u = \frac{u'}{V_0}, t = \frac{t' V_0^2}{4\nu}, \omega = \frac{4\omega' \nu}{V_0^2}, Q_2 = \frac{\nu^2 (C' - C'_\infty) Q_r}{k V_0^2 (T'_w - T'_\infty)}, \\
 \theta &= \frac{T' - T'_\infty}{T'_w - T'_\infty}, \phi = \frac{C' - C'_\infty}{C'_w - C'_\infty}, K = \frac{K' V_0^2}{\nu^2}, \text{Pr} = \frac{\nu \rho c_p}{k}, h = \frac{V_0 L'}{\nu}, Kr = \frac{K_r \nu}{V_0^2}, Sc = \frac{\nu}{D}, \\
 M &= \frac{\sigma B_0^2 \nu}{\rho V_0^2}, Gr = \frac{\nu \beta g (T'_w - T'_\infty)}{V_0^3}, Gc = \frac{\nu \beta^* g (C'_w - C'_\infty)}{V_0^3} \quad \dots(7)
 \end{aligned}$$

In perspective of Equations (6) and (7), Equations (2), (3) and (4) can be lessened to the accompanying dimensionless structure.

$$\frac{1}{4} \frac{\partial u}{\partial t} - (1 + \varepsilon A e^{i\omega t}) \frac{\partial u}{\partial y} = \frac{\partial^2 u}{\partial y^2} + Gr\theta + Gc\phi - \left(M + \frac{1}{K} \right) u \quad \dots(8)$$

$$\frac{1}{4} \frac{\partial \theta}{\partial t} - (1 + \varepsilon A e^{i\omega t}) \frac{\partial \theta}{\partial y} = \frac{1}{\text{Pr}} \left(Q_2 \phi + \frac{\partial^2 \theta}{\partial y^2} \right) \quad \dots(9)$$

$$\frac{1}{4} \frac{\partial \phi}{\partial t} - (1 + \varepsilon A e^{i\omega t}) \frac{\partial \phi}{\partial y} = \frac{1}{Sc} \frac{\partial^2 \phi}{\partial y^2} - Kr\phi \quad \dots(10)$$

The corresponding dimensionless boundary conditions are –

$$\begin{aligned}
 u &= h \frac{\partial u}{\partial y}, \quad \theta = 1 + \varepsilon e^{i\omega t}, \quad \phi = 1 + \varepsilon e^{i\omega t} \quad \text{at} \quad y = 0 \\
 u &\rightarrow 0, \quad \theta \rightarrow 0, \quad \phi \rightarrow 0 \quad \text{as} \quad y \rightarrow \infty \quad \dots(11)
 \end{aligned}$$

Solution of problem

The conditions (8) to (10) are coupled, non-straight fractional differential conditions and these can't be fathomed in shut structure. Be that as it may, these conditions can be decreased to an arrangement of common differential conditions, which can be settled

diagnostically. So this should be possible, when the sufficiency of motions is little, we can expect the arrangements of stream speed temperature field and focus in the area of the plate as:

$$u(y, t) = u_0(y) + \varepsilon e^{i\omega t} u_1(y) + O(\varepsilon^2) + \dots \quad \dots(12)$$

$$\theta(y, t) = \theta_0(y) + \varepsilon e^{i\omega t} \theta_1(y) + O(\varepsilon^2) + \dots \quad \dots(13)$$

$$\phi(y, t) = \phi_0(y) + \varepsilon e^{i\omega t} \phi_1(y) + O(\varepsilon^2) + \dots \quad \dots(14)$$

Substituting (12), (13) and (14) in Equations (8) - (10) and comparing consonant and non-symphonious terms, and dismissing the higher request terms of $O(\varepsilon^2)$, we get –

$$u_0'' + u_0' - F_5 u_0 = -[Gr\theta_0 + Gc\phi_0] \quad \dots(15)$$

$$u_1'' + u_1' - F_9 u_1 = -[Gr\theta_1 + Gc\phi_1 + Au_0'] \quad \dots(16)$$

$$\theta_0'' + Pr\theta_0' = -PrQ_2\phi_0 \quad \dots(17)$$

$$\theta_1'' + Pr\theta_1' - PrG_3\theta_1 = -Pr\phi_1Q_2 - APr\theta_0' \quad \dots(18)$$

$$\phi_0'' + Sc\phi_0' - ScKr\phi_0 = 0 \quad \dots(19)$$

$$\phi_1'' + Sc\phi_1' - \left(\frac{i\omega}{4} + Kr\right)Sc\phi_1 = -ASc\phi_0' \quad \dots(20)$$

where the primes denote the differentiation with respect to y .

The corresponding boundary conditions can be written as

$$\begin{aligned} u_0 = h\left(\frac{\partial u_0}{\partial y}\right), \quad u_1 = h\left(\frac{\partial u_1}{\partial y}\right), \quad \theta_0 = 1, \quad \theta_1 = 1, \quad \phi_0 = 1, \quad \phi_1 = 1 \quad \text{at} \quad y = 0 \\ u_0 \rightarrow 0, \quad u_1 \rightarrow 0, \quad \theta_0 \rightarrow 0, \quad \theta_1 \rightarrow 0, \quad \phi_0 \rightarrow 0, \quad \phi_1 \rightarrow 0 \quad \text{as} \quad y \rightarrow \infty \dots(21) \end{aligned}$$

The legitimate courses of action of conditions (15) – (20) with satisfying the farthest point conditions (21) are given by –

$$u_0 = F_8 e^{-m_7 y} + F_6 e^{-Pr y} + F_7 e^{-m_2 y} \quad \dots(22)$$

$$u_1 = F_{13} e^{-m_3 y} + F_{15} e^{-m_9 y} + F_{10} e^{-m_7 y} + F_{11} e^{-Pr y} + F_{12} e^{-m_4 y} + F_{14} e^{-m_2 y} \quad \dots(23)$$

$$\theta_0 = G_1 e^{-m_2 y} + G_2 e^{-Pr y} \quad \dots(24)$$

$$\theta_1 = F_4 e^{-m_4 y} + F_1 e^{-Pr y} + F_2 e^{-m_2 y} + F_3 e^{-m_3 y} \quad \dots(25)$$

$$\phi_0 = e^{-m_2 y} \quad \dots(26)$$

$$\phi_1 = N_5 e^{-m_3 y} + N_4 e^{-m_2 y} \quad \dots(27)$$

In perspective of the above arrangements, the speed, temperature and concentration appropriations in the limit layer gets to be

$$u(y, t) = u_0 + \varepsilon e^{i\omega t} u_1 = (F_8 e^{-m_7 y} + F_6 e^{-Pr y} + F_7 e^{-m_2 y}) + \varepsilon e^{i\omega t} (F_{10} e^{-m_7 y} + F_{11} e^{-Pr y} + F_{12} e^{-m_4 y} + F_{13} e^{-m_4 y} + F_{14} e^{-m_3 y} + F_{15} e^{-m_9 y}) \quad \dots(28)$$

$$\theta(y, t) = \theta_0 + \varepsilon e^{i\omega t} \theta_1 = (G_2 e^{-Pr y} + G_1 e^{-m_2 y}) + \varepsilon e^{i\omega t} (F_4 e^{-m_4 y} + F_1 e^{-Pr y} + F_2 e^{-m_2 y} + F_3 e^{-m_3 y}) \quad \dots(29)$$

$$\phi(y, t) = \phi_0 + \varepsilon e^{i\omega t} \phi_1 = e^{-m_2 y} + \varepsilon e^{i\omega t} (N_4 e^{-m_2 y} + N_5 e^{-m_3 y}) \quad \dots(30)$$

It is presently essential to ascertain the physical amounts of essential interest, which are the neighborhood divider shear push, the nearby surface warmth, and mass flux. All these are appeared through charts without giving numerical examination. The constants are not offered because of purpose of curtness.

Skinfriction coefficient or shearing stress –

$$\begin{aligned} &= - \left(\frac{\partial u}{\partial y} \right)_{y=0} \\ &= (m_7 F_8 + Pr F_6 + m_7 F_7) + \\ &\varepsilon e^{i\omega t} (m_7 F_{10} + Pr F_{11} + m_4 F_{12} + m_4 F_{13} + m_3 F_{14} + m_9 F_{15}) \quad \dots(31) \end{aligned}$$

Nusslet number

$$\begin{aligned}
&= - \left(\frac{\partial \theta}{\partial y} \right)_{y=0} \\
&= (\text{Pr } G_2 + m_2 G_1) + \varepsilon e^{i\omega t} (m_4 F_4 + \text{Pr } F_1 + m_2 F_2 + m_3 F_3) \quad \dots(32)
\end{aligned}$$

Sherwood number

$$\begin{aligned}
&= - \left(\frac{\partial c}{\partial y} \right)_{y=0} \\
&= m_2 + \varepsilon e^{i\omega t} (m_2 N_4 + m_3 N_5) \quad \dots(33)
\end{aligned}$$

RESULTS AND DISCUSSION

The definition of the issue that records for the impact of radiation assimilation and synthetic response on transient MHD free convective stream over a vertical plate through permeable media was expert in the previous segments. As an aftereffect of the numerical figurings, the dimensionless speed, temperature and focus appropriations for the stream under thought are acquired and their conduct have been talked about for varieties in the overseeing parameters viz., the warm Grashof number Gr, changed Grashof number Gc, attractive field parameter M, retention radiation parameter Q_1 , Prandtl number Pr, and Schmidt number Sc. In the present study, the accompanying default parametric qualities are embraced.

For plotting diagrams the accompanying qualities were embraced: Sc = 0.60, Pr = 0.71, K = 0.5, Gr = 5, Gc = 5, M = 1, Kr = 0.5, A = 0.5, $\omega = 0.01$, h = 0.1, $Q_2 = 1$. All charts hence compare to these unless particularly demonstrated on the fitting diagram. Fig. 1 demonstrates the conduct of fixation for various estimations of synthetic response parameter Kr. It is watched that an expansion in prompts a diminishing in the estimations of focus. Fig. 2 demonstrates the conduct of Temperature for various estimations of substance response parameter Kr. It is watched that an expansion in prompts an abatement in the estimations of Temperature. Fig. 3 demonstrates the conduct of velocity for various estimations of substance response parameter Kr. It is watched that an expansion in prompts a lessening in the estimations of Velocity. The impact of the Schmidt number on the focus is appeared in Fig. 4. As the Sc builds the focus diminishes. The impact of the Schmidt number on the temperature is appeared Fig. 5. As the Sc expands the temperature diminishes. The impact of the Schmidt number on the velocity is appeared Fig. 6. As the Sc expands the velocity increments. The impact of warmth retention coefficient Q_2 on the temperature is appeared Fig. 7. The impact of warmth retention coefficient Q_2 on the velocity is appeared

Fig. 8. Physically the nearness of warmth absorption (thermal sink) impact tends to diminish the liquid temperature. This causes the warm lightness impacts to diminish bringing about a net decrease in the liquid speed. The conduct is plainly clear from Fig. 7 in which the temperature dispersion diminish as Q_2 increments. Additionally the conduct of the speed diminishes as Q_2 increments. It is likewise watched that warm limit layer diminish as the warmth retention impact increment and hydrodynamic increments as the warmth assimilation impact increments in Fig. 8. The merged impact of substance response parameter concerning Schmidt number over Sherwood is seen Fig. 9. It is watched that, expansion in K_r adds to increment of sherwood number. The united impact of warmth retention parameter as for Prandtal number over Nusselt number is seen Fig. 10. It is watched that, expansion in Q_2 adds to diminishing of Nusselt number.

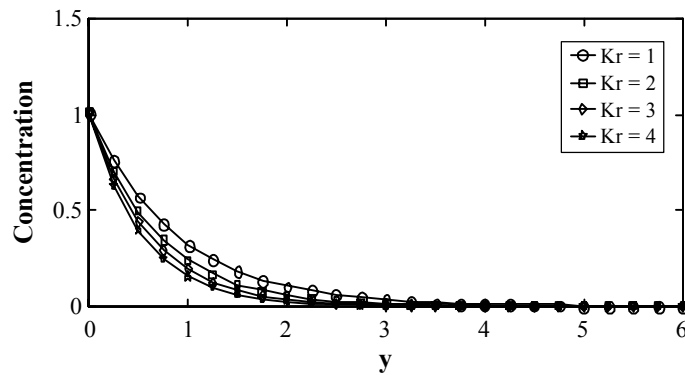


Fig. 1: Impacts of K_r on concentration profiles with $t = \pi/2$, $Sc = 0.60$, $A = 0.5$, $\epsilon = 0.01$

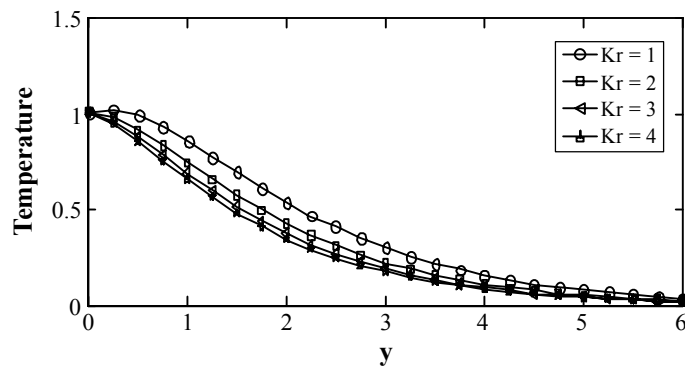


Fig. 2: Impacts of K_r on temperature profiles with $t = \pi/2$, $Sc = 0.60$, $A = 0.5$, $\epsilon = 0.01$, $Pr = 0.71$, $Q_2 = 1$, $\omega = 0.01$

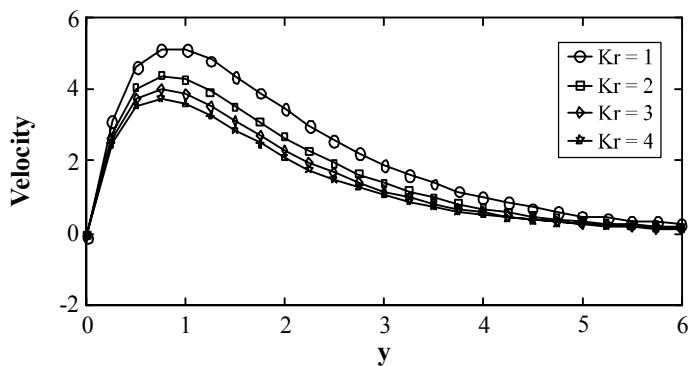


Fig. 3: Impacts of Kr on velocity profiles with $t = \pi/2$, $Sc = 0.60$, $A = 0.5$, $\varepsilon = 0.01$, $Pr = 0.71$, $Q_2 = 1$, $\omega = 0.001$, $M = 1$, $Gr = 20$, $Gc = 10$, $h = 0.01$

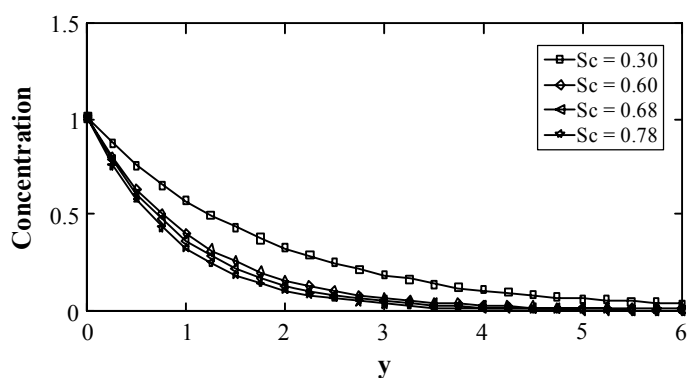


Fig. 4: Effects of Sc on concentration profiles with $t = \pi/2$, $Kr = 0.5$, $A = 0.5$, $\varepsilon = 0.01$

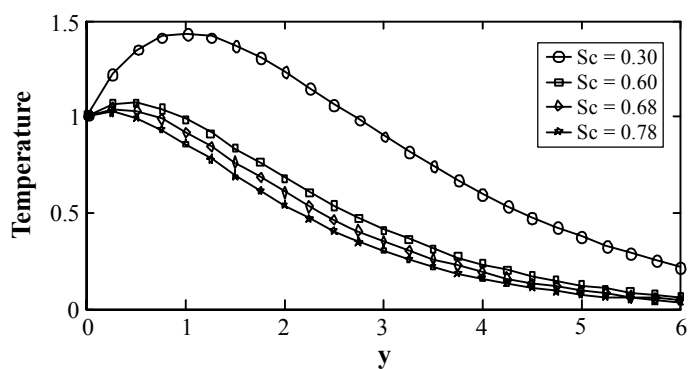


Fig. 5: Impacts of Sc on Temperature profiles with $t = \pi/2$, $A = 0.5$, $\varepsilon = 0.01$, $Pr = 0.71$, $Q_2 = 1$, $Kr = 0.5$, $\omega = 0.01$

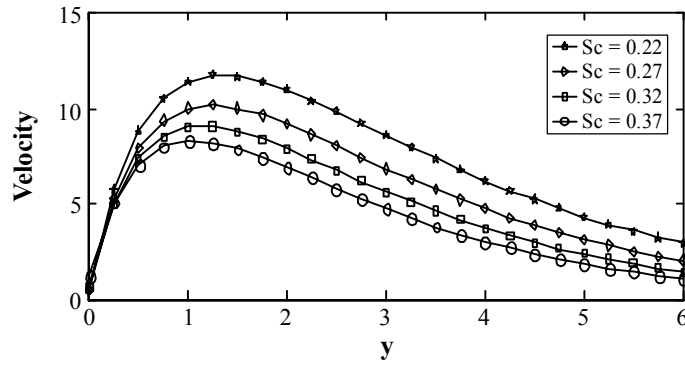


Fig. 6: Impacts of Sc on velocity profiles with $t = \pi/2$, $A = 0.5$, $\omega = 0.01$, $Pr = 0.71$, $Q_2 = 1$, $Kr = 0.5$, $\epsilon = 0.001$, $M = 1$, $Gr = 20$, $Gc = 10$, $h = 0.01$

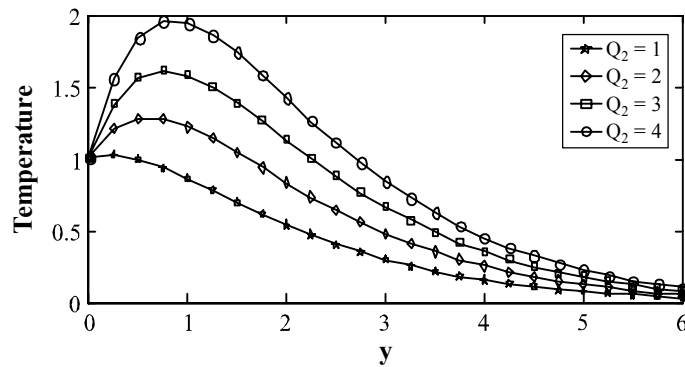


Fig. 7: Impacts of Q_2 on Temperature profiles with $t = \pi/2$, $A = 0.5$, $\epsilon = 0.01$, $Pr = 0.71$, $Sc = 0.60$, $Kr = 0.5$, $\omega = 0.01$

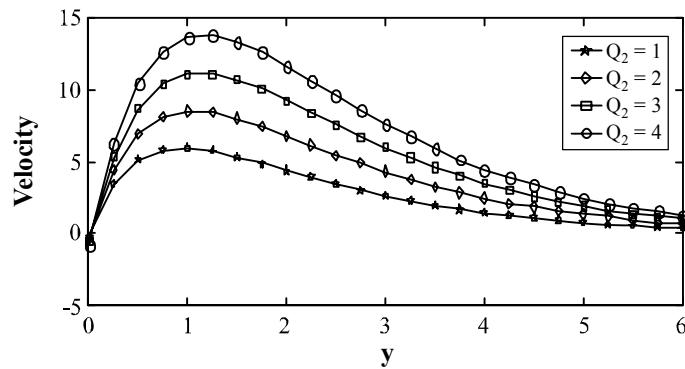


Fig. 8: Impacts of Q_2 on velocity profiles with $t = \pi/2$, $A = 0.5$, $\epsilon = 0.01$, $Pr = 0.71$, $Kr = 0.5$, $\omega = 0.001$, $M = 1$, $Gr = 20$, $Gc = 10$, $h = 0.01$

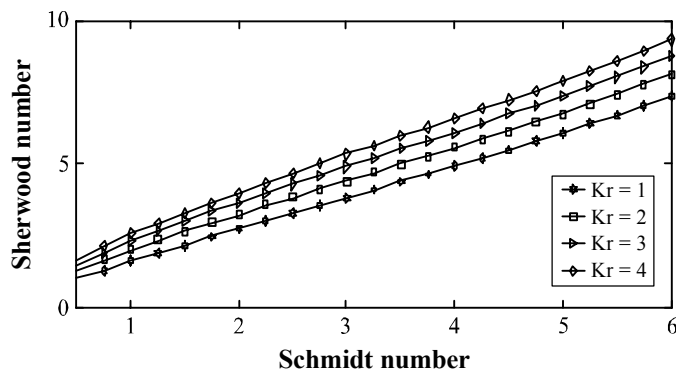


Fig. 9: Impacts of Kr on Sherwood number with $t = \pi/2$; $\omega = 0.1$; $A = 0.5$; $\varepsilon = 0.001$

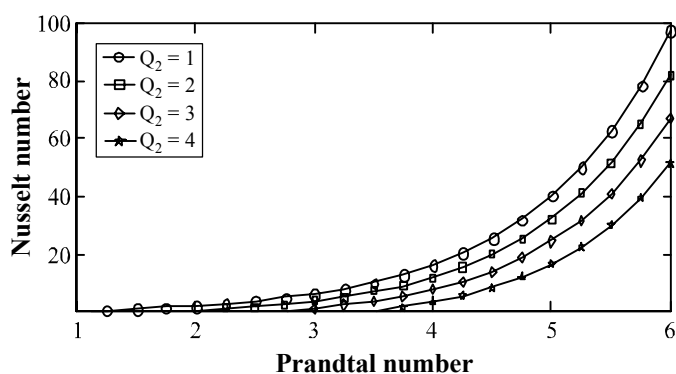


Fig. 10: Impacts of Q₂ on Nusselt number number with Sc = 0.60; $t = \pi/2$; kr = 0.5; $\omega = 0.1$; $A = 0.5$; $\varepsilon = 0.1$

REFERENCES

1. N. G. Kafousias and G. A. Georgantopoulos, Magnetohydrodynamic Free Convection Effects on the Stokes Problem for an Incompressible Viscous Fluid Past an Infinite Vertical Limiting Surface, *Astrophysics and Space Science*, **85(1-2)**, 297-307 (1982).
2. A. Raptis, and V. M. Soundalgekar, MHD Flow Past a Steadily Moving Infinite Vertical Porous Plate with Mass Transfer and Constant Heat Flux, *ZAMM – J. Appl. Mathe. Mech.*, **64(2)**, 127-130 (1984).
3. A. A. Raptis and V. M. Soundalgekar, MHD Flow Past a Steadily Moving Infinite Vertical Porous Plate with Constant Heat Flux, *Nuclear Engineering and Design*, **72(3)**, 373-379 (1982).

4. W. G. England and A. F. Emery, Thermal Radiation Effects on the Laminar Free Convection Boundary Layer of an Absorbing Gas, *J. Heat Transfer*, **91**, 37-44 (1969).
5. A. Raptis and C. V. Massalas, Magnetohydrodynamic Flow Past a Plate by the Presence of Radiation, *Heat and Mass Transfer*, **34**, 107-109 (1998).
6. M. A. Mansour, N. F. El-Anssary and A. M. Aly, Effects of Chemical Reaction and Thermal Stratification on MHD Free Convective Heat and Mass Transfer Over a Vertical Stretching Surface Embedded in a Porous Media Considering Soret and Dufour Numbers, *J. Chem. Engg.*, **145(2)**, 340-345 (2008).
7. A. S. Gupta, Steady and Transient Free Convection of an Electrically Conducting Fluid from a Vertical plate in the Presence of a Magnetic Field, *Appl. Sci. Res.*, **9(1)**, 319-333 (1960).
8. A. J. Chamkha, Transient Hydromagnetic Three-Dimensional Natural Convection from an Inclined Stretching Permeable Surface, *Chem. Engg. J.*, **76(2)**, 159-168 (2000).
9. M. A. Hossain, D. A. S. Rees and I. Pop, Free Convection-Radiation Interaction from an Isothermal Plate Inclined at a Small Angle to the Horizontal, *Acta Mechanica*, **127(1-4)**, 63-73 (1998).
10. S. Shateyi, Thermal Radiation and Buoyancy Effects on Heat and Mass Transfer Over a Semi-Infinite Stretching Surface with Suction and Blowing, *J. Appl. Mathe.*, **12**, Article ID 414830 (2008).
11. S. Shateyi and S. S. Motsa, Thermal Radiation Effects on Heat and Mass Transfer Over an Unsteady Stretching Surface, *Math Problems in Engg.*, **13**, Article ID 965603 (2009).
12. V. Aliakbar, A. Alizadeh-Pahlavan and K. Sadeghy, The Influence of Thermal Radiation on MHD Flow of Maxwellian Fluids Above Stretching Sheets, *Commun Nonlinear Science Numerical Simulations*, **14(3)**, 779-794 (2009).
13. T. Hayat and M. Qasim, Influence of Thermal Radiation and Joule Heating on MHD Flow of a Maxwell Fluid in the Presence of Thermophoresis, *Int. J. Heat Mass Transfer*, **53**, 4780-4791 (2010).
14. R. Cortell, Suction, Viscous Dissipation and Thermal Radiation Effects on the Flow and Heat Transfer of a Power-Law Fluid Past an Infinite Porous Plate, *J. Chem. Engg.*, **89**, 85-93 (2011).

15. J. L. Ramprasad, K. S. Balamurugan and G. Dharmiah, Unsteady MHD Convective Heat And Mass Transfer Past An Inclined Moving Surface With Heat Absorption, JP J. Heat Mass Transfer., **13(1)**, 33-51 (2016).
16. G. V. Ramana Reddy, N. Bhaskar Reddy and A. J. Chamkha, Mhd Mixed Convection Oscillatory Over a Vertical Surface in a Porous Medium with Chemical Reaction and Thermal Radiation, J. Appl. Fluid Mech., **9(3)**, 1221-1229 (2016).
17. P. Singh and C. B. Gupta, Mhd Free Convective Flow of Viscous Fluid Through a Porous Medium Bounded by an Oscillating Porous Plate in Slip Flow Regime with Mass Transfer, Indian J. Theoretical Physics, **52(2)**, 111-120 (2005).
18. S. D. Nigam and S. N. Singh, Heat Transfer by Laminar Flow Between Parallel Plates Under the Action of Transverse Magnetic Field, Q. J. mach. Appl. Math., **13**, 85-92 (1960).
19. D. R. V. P. Rao, D. V. Krishna and L. Debnath, Combined Effect of Free and Forced Convection on Mhd Flow in a Rotating Porous Channel, Acta Mech., **34**, 225-240 (1982).
20. G. Sarojamma and D. V. Krishna, Transient Hydromagnetic Convective Flow in a Rotating Channel with Porous Boundaries, Acta Mech., **40**, 277-288 (1981).
21. R. Siva Prasad, Convection Flows in Magneto Hydro Dynamics, Ph.D., Thesis, S. K. University, Anantapur, India (1985).

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