



## EVALUATION OF TRANSITION TEMPERATURE $T_C$ AND ISOTOPE EFFECT 'A' IN ALKALI DOPED FULLERIDES

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### ABSTRACT

Using the theoretical formalism of Kresin (1991), which includes contribution for both intermolecular and intramolecular phonon coupling, the transition temperature  $T_C$ , isotope effect exponent  $\alpha$  for  $K_3C_{60}$  fulleride was evaluated and theoretical result are in good agreement with the experimental data and also with other theoretical worker's.

**Key words:** Intermolecular phonon mod, Intramolecular phonon mode, Transition temperature  $T_C$ , Isotope effect exponent  $\alpha$ , Energy gap parameter.

### INTRODUCTION

The fullerenes ( $C_{60}$ ,  $C_{70}$ ...) have attracted much interest since their discovery<sup>1</sup>. The interest increased dramatically when it was discovered<sup>2</sup> how to produce  $C_{60}$  in large quantities to make solids (fullerites) of a size that allowed traditional solid state experiment. It was also found that intercalation of alkali metal atom in solid  $C_{60}$  leads to metallic behaviour. It was also found that some of these alkali doped  $C_{60}$  compounds (fullerides) are superconducting with transition temperature  $T_C$  around 30-40 K.  $RbCs_2C_{60}$  has  $T_C = 33$  K while  $Cs_3C_{60}$  under pressure has  $T_C = 40$  K. The great interest of these alkali-doped  $C_{60}$  is the question whether or not such a large value of  $T_C$  can be caused by coupling to phonons alone. There has been a great effort to characterize and understand both the normal state and the superconducting properties of fullriders<sup>2-4</sup>.

$C_{60}$  is a most symmetric molecule. The 60 carbon atoms are all equivalent and form 12 pentagons and 20 hexagons.  $C_{60}$  molecule condensed into solid of weakly bond molecule. For an undoped  $C_{60}$  solid the  $h_u$  band is full and  $t_{1u}$  band is empty and this system is

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therefore band insulator. When solid  $C_{60}$  is doped by alkali atoms, the alkali atoms donate about one electron each of the  $t_{1u}$  band<sup>5-7</sup>.

$K_3C_{60}$  is the subject of renewed interest because of its technological and industrial importance. In close relation the nature of the attracting mechanism with conventional phonon mediated electron pairing unconventional electron mechanism in fullerenes ( $T_C$  -20-45 K) remains a central point of several experimental studies. Neutron-inelastic scattering measurement<sup>8,9</sup> probes that the vibrational spectrum may be conveniently divided into two regions. One of them belongs to the rotation of  $C_{60}$  molecule and the intermolecular vibration ( $\approx 2.5$ -25 meV) and the others are the intramolecular mode with frequency 25-200 meV. Raman scattering<sup>10</sup> yields the on-ball  $C_{60}$  vibrational mode ( $\sim 40$  meV -0.25eV). Lattice dynamical studies point out that the spectrum consists of modes of  $C_{60}$  ( $26\text{ cm}^{-1}$ ), the vibrations of  $K^+$  ions, the translation variation as a whole ( $45$ - $120\text{ cm}^{-1}$ ) and the  $C_{60}$  vibrations ( $\sim 260$ - $1900\text{ cm}^{-1}$ ). The wide frequency range of phonon spectrum raises an important issue to clarify which of these molecular phonon modes induce the superconductivity in doped fullerenes. The energy scales of the various phonon modes that mediate electron-phonon coupling are different. This difference was observed because of the dependence of the transition temperature on isotopic mass. It was also pointed out that it is the alkali isotope effect and not the carbon isotope effect that identifies whether inter or intramolecular phonon modes induce the superconductivity.

In this paper, taking the theoretical formalism of Kresin<sup>11</sup>, we have theoretically evaluated the transition temperature of  $K_3C_{60}$  as a function of intramolecular phonon coupling with various values of Coulomb repulsion parameter  $\mu^*$ . We have also evaluated isotope exponent  $\alpha$  and energy gap parameter  $2\Delta/T_C$ .

### Mathematical formulae used in this

In order to calculate transition temperature  $T_C$  one uses theoretical formalism given by Kresin<sup>11</sup>. This formalism is developed in the strong coupling theory where  $\lambda > 1$ ,  $\lambda$  is coupling strength.

$$T_c^{er} = 0.25 \omega_{er} [\exp(2/\lambda_{eff}) - 1]^{-1/2} \quad \dots(1)$$

$$\text{where } \lambda_{eff} = (\lambda_{er} - \mu^*) [1 + 2\mu^* + \lambda_{er} + \mu^* t(\lambda_{er})]^{-1} \quad \dots(2)$$

$\mu^*$  is the renormalized coulomb repulsive parameter and  $\omega_{er}$  is the characteristic phonon frequency.

$$t(\lambda_{er}) = 1.5 \exp(0.28 \lambda_{er}) \quad \dots(3)$$

$\lambda_{er}$  is the electron-phonon coupling strength. In this formalism both inter or intramolecular phonon modes induces the superconductivity.

One takes<sup>12</sup>

$$\begin{aligned}\lambda_{eff} &= 0.62 \\ t(\lambda_{er}) &= 1.07 \\ \lambda_{er} &= 1.2\end{aligned}\quad \dots(4)$$

Using Eq. (1) it was found that  $T_c^{er} = 5$  K which is much lower than the reported experimental value of  $K_3C_{60}$  (20K). Then equation (1) is modified by taking the simultaneous presence of both inter and intramolecular phonon in fullerenes.

$$T_c = T_c^{eff} [\omega_{ra} / T_c^{eff}]^\chi \quad \dots(5)$$

$$\text{Where, } \chi = \lambda_{ra} [\lambda_{ra} + \lambda_{er}]^{-1} \quad \dots(6)$$

$\lambda_{ra}$  is the intramolecular coupling constant and is obtained from the scattering time<sup>13</sup>.

$$\lambda_{ra} = \hbar / (2\pi K_\beta T \tau) \quad \dots(7)$$

$\tau$  is the scattering time,  $[\tau = 1.2 \times 10^{-14} \text{s}]^{13}$ ,  $\omega_{ra} = 1012 \text{ cm}^{-1}$

### Isotope effect coefficient

The isotope effect coefficient is defined as -

$$\alpha = (0.5) (d \ln T_c / d \ln m_{red}) \quad \dots(8)$$

$m_{red}$  is reduced molecular mass. The  $m_{red}$  dependence of  $T_c$  is introduced in terms of  $\mu^*$ , then the value of  $\alpha$  becomes-

$$\alpha = \{1 - \mu^{*2} (1 + 2\lambda_{er} + \lambda_{er}^2) / \lambda_{eff} (\lambda_{er} - \mu^*) [1 + (4T_c / \omega)^2]\} / 2 \quad \dots(9)$$

Here  $\omega$  is the average phonon frequency for  $K_3C_{60}$  Fullerene ( $T_c = 20$ ) one uses

$$\lambda_{er} = 1.2 \text{ and } \lambda_{ra} = 0.4$$

Energy gap parameter-

The energy gap parameter<sup>11</sup>  $\beta$  is given by -

$$\begin{aligned}B &\approx 2\Delta(0) / H_s T_c \\ &\approx 3.5 [1 + 5 (T_c / \omega_{ra})^2 \ln(\omega_{ra} / T_c)]\end{aligned}\quad \dots(10)$$

Screening parameter

The renormalized coulomb repulsion parameter is defined as -

$$\mu^* = [(1 + \mu) \ln(\epsilon_F / \omega_{\text{eff}})] \quad \dots(11)$$

The coulomb strength parameter  $\mu$  is given by -

$$\mu = N(\epsilon_F)U \quad \dots(12)$$

Where  $N(\epsilon_F)$  is the density of states at the Fermi energy and  $U$  is the static screened interaction,  $V(q, \omega=0)$  average over the Fermi sphere. One calculate  $\mu$  from static dielectric function.

$\epsilon(q)$  in the long wavelength limit.

$$\epsilon(q) = (4m_e^2 K_F) [\pi \hbar^2 q^2]^{-1} \quad \dots(13)$$

Using Eq. (13),  $\mu$  is calculated as –

$$\mu = 1/2 \pi a_B k_B \ln[1 + 3 \pi a_B k_F / 2] \quad \dots(14)$$

$a_B$  = Bohr radius

$K_F$  = Fermi wave vector

## RESULTS AND DISCUSSION

In this paper, we have presented a method of evaluation of Superconducting transition temperature  $T_c$  of alkali metal doped fullerenes. We have taken the theoretical formalism developed by Kresin<sup>11</sup> to evaluate  $T_c$ .  $T_c$  was evaluated as function of intramolecular coupling strength  $\lambda_{\text{ra}}$ . The result is shown in Table 2. It is clear from the result that  $T_c$  is strongly influenced by the Coulomb repulsive parameter  $\mu^*$ .  $T_c$  is higher for smaller values of  $\mu^*$ . The higher value of the  $\mu^*$  is the result of the increased electron-phonon repulsive contribution alongwith the intramolecular phonons producing an attractive interaction. This attractive interaction is reduced alone against the phonon attraction and attributes to suppress  $T_c$ . Higher and positive  $\mu^*$  implies a constant repulsive interaction that is insufficient to create a superconductivity state despite of coupling of intramolecular phonons. This shows that the Coulombs pseudopotential parameter is an important concept in fullerenes. It appear from Table 2 that  $T_c$  is strongly sensitive to  $\lambda_{\text{ra}}$  even for moderate coupling. One can see that for fixed value of  $\lambda_{\text{ra}} = 0.35$  one gets  $T_c = 30.48$  K for  $\mu^* = 0.0$ ,  $T_c = 23.27$  K for  $\mu^* = 0.1$  and  $T_c = 19.29$  for  $\mu^* = 0.2$ . Thus, it appears that if one start with a pure intramolecular

phonon mechanism and later on adding a correction term with moderately coupled a intramolecular phonons. One can easily enhance  $T_c$  value. In table 1, we show the transition temperature  $T_c$  of some alkali metal doped  $C_{60}$  compound using equation (9), we have evaluated the isotope effect exponent ' $\alpha$ ' for  $K_3C_{60}$  as a function of  $\mu^*$  keeping the fixed value of  $\lambda_{er} = 1.2$  and  $\lambda_{ra} = 0.4$ . The results are shown in Table 3. Our theoretically evaluated results show that isotope exponent  $\alpha$  decreases with increase value of  $\mu^*$ . Our theoretically obtained result is in agreement with the experimental data<sup>14</sup> and also with other theoretical workers<sup>15,16</sup>. In table 4, we have given the values of energy gap parameter  $2\Delta/T_c$  obtained with various experimental studies<sup>17-26</sup> for  $Rb_6C_{60}$  and  $K_3C_{60}$  compounds.

**Table 1: Transition temperature  $T_c$  of some alkali metal doped  $C_{60}$  compounds**

Compounds	$T_c$ (K)
$K_3C_{60}$	20
$K_2RbC_{60}$	22
$Rb_2KC_{60}$	25
$Rb_3C_{60}$	29
$Rb_2CsC_{60}$	31
$CsRbC_{60}$	33
$Cs_3C_{60}$	47
$Rb_{2.7}Tl_{2.2}C_{60}$	45

**Table 2: An evaluated result of superconducting transition temperature  $T_c$  as a function of  $\lambda_{ra}$  (intramolecular coupling constant) for different values of coulomb repulsive parameter  $\mu^*$ . The results are compared with alkali doped  $C_{60}$  ( $K_3C_{60}$ )**

$\lambda_{er}$	Theory $T_c$ (K)				Expt. $T_c$ (K)
	$\mu^* = 0.0$	$\mu^* = 0.1$	$\mu^* = 0.2$	$\mu^* = 0.3$	
0.00	11.25	6.78	4.78	3.25	20.0
0.10	13.68	9.49	7.16	5.10	20.0
0.20	16.38	12.58	10.24	6.28	20.0
0.25	21.59	16.16	12.59	8.46	20.0
0.30	24.26	18.25	15.56	10.29	20.0

Cont...

$\lambda_{\text{er}}$	Theory $T_c$ (K)				Expt. $T_c$ (K)
	$\mu^* = 0.0$	$\mu^* = 0.1$	$\mu^* = 0.2$	$\mu^* = 0.3$	
0.35	30.48	20.16	17.18	12.50	20.0
0.40	35.12	23.27	19.29	14.56	20.0
0.45	38.46	27.18	21.86	16.29	20.0
0.50	42.56	31.56	24.28	18.58	20.0
0.55	48.29	33.29	27.47	22.25	20.0
0.60	52.18	37.54	31.26	25.18	20.0
0.65	58.29	43.56	34.18	29.56	20.0
0.70	62.17	48.28	38.56	31.48	20.0
0.75	67.29	54.62	40.14	35.29	20.0
0.80	70.18	59.10	44.25	40.15	20.0

**Table 3: An evaluated result of isotope effect exponent  $\alpha$  as a function of Coulomb pseudopotential parameter  $\mu^*$  for  $\text{K}_3\text{C}_{60}$  Fulleride,  $\lambda_{\text{er}} = 1.2$ ,  $\lambda_{\text{ra}} = 0.4$**

$\mu^*$	Isotope exponent $\alpha$
0.00	0.525
0.025	0.517
0.050	0.502
0.075	0.476
0.10	0.457
0.12	0.405
0.14	0.386
0.15	0.325
0.16	0.305
0.18	0.269
0.20	0.248
0.22	0.207
0.24	0.196
0.25	0.108

**Table 4: An experimental result of energy gap parameter  $2\Delta/T_C$  for various alkali doped fulloride**

Fulleride	$2\Delta/T_C$	References	
Rb <sub>6</sub> C <sub>60</sub>	5.3	Z. Zang et al.	1991
	5.4	Jess et al.	1994
	2.0-4.0	Jess et al.	1996
	4.3	Tycho et al.	1992
	3.6	Kiefl et al.	1993
	3.48	Degorgi et al.	1995
	4.1	Gu et al.	1994
K <sub>3</sub> C <sub>60</sub>	3.0	Tycko et al.	1992
	4.7	Sasaki et al.	1994
	3.4 ± 0.4	Ausen-Sinzie et al.	1993
	3.44	Degiorgi et al.	1994
$(2\Delta/T_C)_{BCS} = 4.2$			

## REFERENCES

1. H. W. Kroto, J. R. Heath, S. C. O'Brien, R. F. Curl and R. E. Smalley, *Nature*, **318**, 162 (1985).
2. M. G. Mitch, S. J. Chese and J. S. Lannien, *Phys. Rev. Lett. (PRL)*, **68**, 883 (1992).
3. M. G. Mitch and J. S. Lannien, *Phys. Rev. B*, **48**, 16192 (1993).
4. V. N. Kostur and B. Matrovic, *Phys. Rev. B*, **50**, 12774 (1994).
5. A. S. Alexandov and V. V. Kabanov, *Phys. Rev. B*, **54**, 3655 (1996).
6. E. Cappelluti, C. Grimaldi, L. Piefronero and S. Strassler, *Phys. Rev. Lett. (PRL)*, **85**, 4771 (2000).
7. T. Chida, S. Suzuki and K. Nakao, *J. Phys. Soc. Jpn.*, **71**, 525 (2002).
8. K. Prasseney, J. Tomkinson, C. Christedis, M. J. Rossensky, D. W. Marphy and R. C. Hadden, *Nature*, **354**, 462 (1991).
9. L. Pintschovices, *Rep. Pro. Phys.*, **59**, 473 (1996).

10. L. Ferro and L. Mihaly, *Rep. Prog. Phys.*, **64**, 649 (2001).
11. V. Z. Kresin, *Phy. Rev. B*, **46**, 14833 (1991).
12. M. L. Zhang and H. Y. Guo, *Physica C*, 227 (1994).
13. X. D. Xiary, J. G. Hau, V. H. Crespi, A. Zettl and M. L. Cohen, *Nature*, **361**, 54 (1993).
14. B. K. Agrawal, P. S. Yadav, P. Srivastava, R. Srivastava and S. Singh, *J. Superconducting: In Incorporating Novel Magnetism*, **17**, 289 (2004).
15. D. Varshney, A. Dube, K. K. Choudhary and R. K. Singh, *Bull. Mater. Sci.*, **28**, 155 (2005).
16. D. Varshney, *J. Superconductivity*, **13**, 171 (2000).
17. Z. Zhang, C. C. Chen and C. M. Lieber, *Science*, **254**, 1619 (1991).
18. P. Jess, S. Behler, M. Bernaconi and H. J. Guntherodt, *Physica C*, 235-240, **2499** (1994).
19. P. Jess, U. Hubler, S. Behler and H. J. Guntherodt, *Synth. Met.*, **72**, 201 (1996).
20. R. Tycho, G. Dabbagh, D. W. Murphy, Q. Zhu and J. E. Fisher, *Phys. Rev. B*, **48**, 9097 (1993).
21. R. F. Kiefl, W. A. Marfarlam, K. H. Chow, S. Dunsiger, T. L. Duty and J. E. Fisher *Phys. Rev. Lett. (PRL)*, **70**, 3987 (1993).
22. L. Degiorgi, G. Briceno, M. S. Fulther, A. Zttle and P. Wachter, *Nature*, **369**, 541 (1994).
23. L. Degiorgi, *Mod. Phy. Lett. B*, **9**, 445 (1995).
24. C. Gu, B. W. Veal, R. Liu, A. P. Paulikas, H. H. Wang and J. W. Williams, *Phys. Rev. B*, **50**, 16566 (1994).
25. S. Sasaki, A. Matsuda and C. W. Chu, *J. Phys. Soc. Jpn.*, **63**, 167 (1994).
26. Auban-Senzier, P. G. Quirion, D. Jerome, P. Bernier and A. Bassat, *Synth. Met.*, **56**, 3027 (1993).

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